

STABILIZATION OF A GENERALIZED TELEGRAPH EQUATION ON STAR-SHAPED NETWORK

Alaa HAYEK

Directors of Thesis: **Pr. Serge Nicaise** and **Pr. Ali Wehbe**

Co-Director: **Dr. Zaynab Salloum**

UVHC, LAMAV,FR CNRS 2956 , FRANCE.
UNIVERSITÉ LIBANAISE, EDST, EQUIPE EDP-AN,
HADATH, BEYROUTH, LIBAN.

30/10/2019

1 INTRODUCTION

2 RESULTS ACHIEVED

- Stabilization of a generalized Telegraph equation on a one dimensional star shaped network

3 CONTINUATION OF THE WORK

4 REFERENCES

In [1], S. Nicaise considered the stabilization of the generalized telegraph equation set in a real interval:

$$\begin{cases} V_t + gV + aI_x + kW & = 0, & \text{in } (0, L) \times (0, \infty), \\ I_t + rI + bV_x & = 0, & \text{in } (0, L) \times (0, \infty), \\ W_t + cW & = V, & \text{in } (0, L) \times (0, \infty), \end{cases} \quad (1)$$

with the following boundary conditions

$$V(0, t) = V(L, t) = 0, \quad t \in \mathbb{R}_+^*,$$

In [1], S. Nicaise considered the stabilization of the generalized telegraph equation set in a real interval:

$$\begin{cases} V_t + gV + aI_x + kW & = 0, & \text{in } (0, L) \times (0, \infty), \\ I_t + rI + bV_x & = 0, & \text{in } (0, L) \times (0, \infty), \\ W_t + cW & = V, & \text{in } (0, L) \times (0, \infty), \end{cases} \quad (1)$$

with the following boundary conditions

$$V(0, t) = V(L, t) = 0, \quad t \in \mathbb{R}_+^*,$$

and the following initial conditions

$$V(x, 0) = V_0(x), \quad I(x, 0) = I_0(x), \quad W(x, 0) = W_0(x), \quad x \in (0, L),$$

In [1], S. Nicaise considered the stabilization of the generalized telegraph equation set in a real interval:

$$\begin{cases} V_t + gV + aI_x + kW & = 0, & \text{in } (0, L) \times (0, \infty), \\ I_t + rI + bV_x & = 0, & \text{in } (0, L) \times (0, \infty), \\ W_t + cW & = V, & \text{in } (0, L) \times (0, \infty), \end{cases} \quad (1)$$

with the following boundary conditions

$$V(0, t) = V(L, t) = 0, \quad t \in \mathbb{R}_+^*,$$

and the following initial conditions

$$V(x, 0) = V_0(x), \quad I(x, 0) = I_0(x), \quad W(x, 0) = W_0(x), \quad x \in (0, L),$$

where, a, b, c, r, k and g are all non-negative functions in $L^\infty(0, L)$ that verify some assumptions mentioned in [1].

History: Results obtained in [1]

In [1], S.Nicaise was interested in studying the energy decay rate of system (1) on a **cable**.

In [1], S.Nicaise was interested in studying the energy decay rate of system (1) on a cable.

- Under the additional condition that $r + g > 0$ in $(0, L)$, an exponential energy decay rate is established.

In [1], S.Nicaise was interested in studying the energy decay rate of system (1) on a **cable**.

- Under the additional condition that $r + g > 0$ in $(0, L)$, an **exponential energy decay rate** is established.
- Otherwise, a **polynomial energy decay rate of type $\frac{1}{t}$** for smooth initial data is established.

History: Results obtained in [1]

In [1], S.Nicaise was interested in studying the energy decay rate of system (1) on a **cable**.

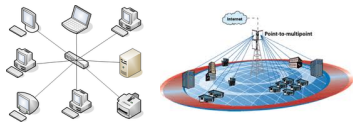
- Under the additional condition that $r + g > 0$ in $(0, L)$, an **exponential energy decay rate** is established.
- Otherwise, a **polynomial energy decay rate of type $\frac{1}{t}$** for smooth initial data is established.

Moreover, the obtained polynomial decay rate is optimal in the case $r = g = 0$.

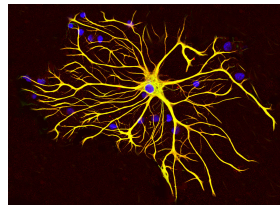


S. Nicaise. Stabilization and asymptotic behavior of a generalized telegraph equation. *Z. Angew. Math. Phys.*, 66(6):3221–3247, 2015.

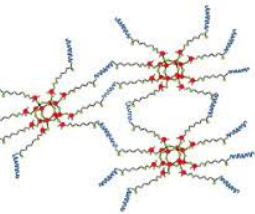
Applications of Networks



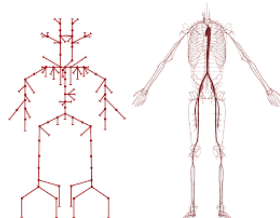
(A) Electrical Networks



© Adapted by A. S. FAR with permission from Macmillan Publishers



(B) Chemical Networks



(C) Biological Networks

1 INTRODUCTION

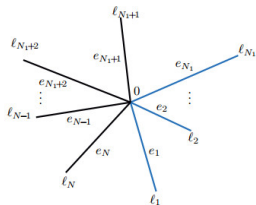
2 RESULTS ACHIEVED

- Stabilization of a generalized Telegraph equation on a one dimensional star shaped network

3 CONTINUATION OF THE WORK

4 REFERENCES

Objective

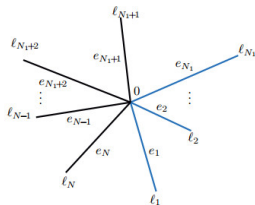


- **Objective:** To study the stabilization of the generalized telegraph equation on a one dimensional **star shaped graph**.

$$\begin{cases} V_t + gV + al_x + kW & = 0, \text{ in } \mathcal{Q} \times (0, \infty), \\ I_t + rI + bV_x & = 0, \text{ in } \mathcal{Q} \times (0, \infty), \\ W_t + cW & = V, \text{ in } \mathcal{Q} \times (0, \infty), \end{cases}$$

with $\mathcal{Q} = \prod_{\ell=0}^N (0, L)$.

Objective



- **Objective:** To study the stabilization of the generalized telegraph equation on a one dimensional **star shaped graph**.

$$\begin{cases} V_t + gV + al_x + kW & = 0, \text{ in } \mathcal{Q} \times (0, \infty), \\ I_t + rI + bV_x & = 0, \text{ in } \mathcal{Q} \times (0, \infty), \\ W_t + cW & = V, \text{ in } \mathcal{Q} \times (0, \infty), \end{cases}$$

with $\mathcal{Q} = \prod_{\ell=0}^N (0, L)$.

- **Main Result:** The energy of the system decays exponentially to zero.

Objective

Model

- Let Γ be a star graph, which is constructed by attaching $N+1$ lines at a common vertex.

Objective

Model

- Let Γ be a star graph, which is constructed by attaching $N+1$ lines at a common vertex.
- We parametrize each edge of Γ by $(0,L)$.

Objective

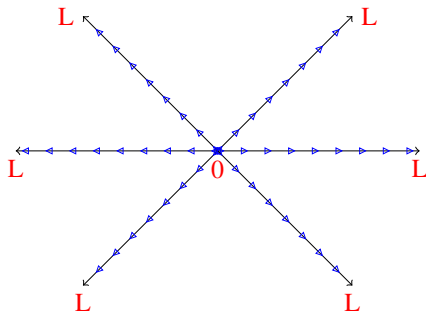
Model

- Let Γ be a star graph, which is constructed by attaching $N+1$ lines at a common vertex.
- We parametrize each edge of Γ by $(0, L)$.
- $\Gamma = \prod_{\ell=0}^N (0, L)$.

Objective

Model

- Let Γ be a star graph, which is constructed by attaching $N+1$ lines at a common vertex.
- We parametrize each edge of Γ by $(0, L)$.
- $\Gamma = \prod_{\ell=0}^N (0, L)$.



Problem

We consider the following generalized telegraph equation on a one-dimensional star shaped network Γ , as follows

$$\begin{cases} V_{\ell,t} + g_{\ell} V_{\ell} + a_{\ell} I_{\ell,x} + k_{\ell} W_{\ell} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ I_{\ell,t} + r_{\ell} I_{\ell} + b_{\ell} V_{\ell,x} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ W_{\ell,t} + c_{\ell} W_{\ell} & = V_{\ell}, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}. \end{cases} \quad (2)$$

This system is a coupling between the usual telegraph equation where the electric unknowns are

Problem

We consider the following generalized telegraph equation on a one-dimensional star shaped network Γ , as follows

$$\begin{cases} V_{\ell,t} + g_{\ell} V_{\ell} + a_{\ell} I_{\ell,x} + k_{\ell} W_{\ell} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ I_{\ell,t} + r_{\ell} I_{\ell} + b_{\ell} V_{\ell,x} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ W_{\ell,t} + c_{\ell} W_{\ell} & = V_{\ell}, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}. \end{cases} \quad (2)$$

This system is a coupling between the usual telegraph equation where the electric unknowns are

- $V = (V_{\ell})_{\ell=0}^N$ (electric potential),

Problem

We consider the following generalized telegraph equation on a one-dimensional star shaped network Γ , as follows

$$\begin{cases} V_{\ell,t} + g_{\ell} V_{\ell} + a_{\ell} I_{\ell,x} + k_{\ell} W_{\ell} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ I_{\ell,t} + r_{\ell} I_{\ell} + b_{\ell} V_{\ell,x} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ W_{\ell,t} + c_{\ell} W_{\ell} & = V_{\ell}, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}. \end{cases} \quad (2)$$

This system is a coupling between the usual telegraph equation where the electric unknowns are

- $V = (V_{\ell})_{\ell=0}^N$ (electric potential),
- $I = (I_{\ell})_{\ell=0}^N$ (electric current),

Problem

We consider the following generalized telegraph equation on a one-dimensional star shaped network Γ , as follows

$$\begin{cases} V_{\ell,t} + g_{\ell} V_{\ell} + a_{\ell} I_{\ell,x} + k_{\ell} W_{\ell} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ I_{\ell,t} + r_{\ell} I_{\ell} + b_{\ell} V_{\ell,x} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ W_{\ell,t} + c_{\ell} W_{\ell} & = V_{\ell}, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}. \end{cases} \quad (2)$$

This system is a coupling between the usual telegraph equation where the electric unknowns are

- $V = (V_{\ell})_{\ell=0}^N$ (electric potential),
- $I = (I_{\ell})_{\ell=0}^N$ (electric current),

and a first order differential equation in an auxiliary variable $W = (W_{\ell})_{\ell=0}^N$ (representing the non-local effects) of parabolic type.

Problem

We consider the following generalized telegraph equation on a one-dimensional star shaped network Γ , as follows

$$\begin{cases} V_{\ell,t} + g_{\ell} V_{\ell} + a_{\ell} I_{\ell,x} + k_{\ell} W_{\ell} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ I_{\ell,t} + r_{\ell} I_{\ell} + b_{\ell} V_{\ell,x} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ W_{\ell,t} + c_{\ell} W_{\ell} & = V_{\ell}, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}. \end{cases} \quad (2)$$

This system is a coupling between the usual telegraph equation where the electric unknowns are

- $V = (V_{\ell})_{\ell=0}^N$ (electric potential),
- $I = (I_{\ell})_{\ell=0}^N$ (electric current),

and a first order differential equation in an auxiliary variable $W = (W_{\ell})_{\ell=0}^N$ (representing the non-local effects) of parabolic type.

All coefficients involved in (2) are in $L^{\infty}(\Omega)$, real valued and non negative. Moreover, we suppose that

Problem

We consider the following generalized telegraph equation on a one-dimensional star shaped network Γ , as follows

$$\begin{cases} V_{\ell,t} + g_{\ell} V_{\ell} + a_{\ell} I_{\ell,x} + k_{\ell} W_{\ell} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ I_{\ell,t} + r_{\ell} I_{\ell} + b_{\ell} V_{\ell,x} & = 0, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}, \\ W_{\ell,t} + c_{\ell} W_{\ell} & = V_{\ell}, \text{ in } (0, L) \times \mathbb{R}_+, \ell \in \{0, \dots, N\}. \end{cases} \quad (2)$$

This system is a coupling between the usual telegraph equation where the electric unknowns are

- $V = (V_{\ell})_{\ell=0}^N$ (electric potential),
- $I = (I_{\ell})_{\ell=0}^N$ (electric current),

and a first order differential equation in an auxiliary variable $W = (W_{\ell})_{\ell=0}^N$ (representing the non-local effects) of parabolic type.

All coefficients involved in (2) are in $L^{\infty}(\Omega)$, real valued and non negative. Moreover, we suppose that

$$a_{\ell} \geq \mu, \quad b_{\ell} \geq \mu, \quad c_{\ell} \geq \mu, \quad k_{\ell} + g_{\ell} \geq \mu \quad \text{a.e in } \Omega, \quad \ell = 0, \dots, N, \quad (3)$$

where μ is a positive real number.

Boundary Conditions

System (2) is subjected to the dissipative boundary condition on all the external vertices

Boundary Conditions

System (2) is subjected to the dissipative boundary condition on all the external vertices

$$V_\ell(L) - \alpha_\ell l_\ell(L) = 0, \quad \alpha_\ell > 0, \quad \forall \ell \in \{0, \dots, N\}, \quad (4)$$

Boundary Conditions

System (2) is subjected to the dissipative boundary condition on all the external vertices

$$V_\ell(L) - \alpha_\ell l_\ell(L) = 0, \quad \alpha_\ell > 0, \quad \forall \ell \in \{0, \dots, N\}, \quad (4)$$

and to the transmission conditions (called the improved Kirchoff conditions) on the interior common vertex

Boundary Conditions

System (2) is subjected to the dissipative boundary condition on all the external vertices

$$V_\ell(L) - \alpha_\ell I_\ell(L) = 0, \quad \alpha_\ell > 0, \quad \forall \ell \in \{0, \dots, N\}, \quad (4)$$

and to the transmission conditions (called the improved Kirchoff conditions) on the interior common vertex

$$\begin{cases} \sum_{k=1}^N Z_{\ell k} I_{k,t}(0) &= \frac{1}{\delta} (V_0(0) - V_\ell(0)), \quad \forall \ell \in \{1, \dots, N\}, \\ V_{0,t}(0) &= \frac{-1}{\delta Y} \sum_{\ell=0}^N I_\ell(0), \end{cases} \quad (5)$$

Boundary Conditions

System (2) is subjected to the dissipative boundary condition on all the external vertices

$$V_\ell(L) - \alpha_\ell I_\ell(L) = 0, \quad \alpha_\ell > 0, \quad \forall \ell \in \{0, \dots, N\}, \quad (4)$$

and to the transmission conditions (called the improved Kirchoff conditions) on the interior common vertex

$$\begin{cases} \sum_{k=1}^N Z_{\ell k} I_{k,t}(0) &= \frac{1}{\delta} (V_0(0) - V_\ell(0)), \quad \forall \ell \in \{1, \dots, N\}, \\ V_{0,t}(0) &= \frac{-1}{\delta Y} \sum_{\ell=0}^N I_\ell(0), \end{cases} \quad (5)$$

where, $Y > 0$, $\delta > 0$, and $Z = (Z_{\ell k})_{N \times N}$ is a symmetric, positive definite matrix.

We set

We set

- $\nu(t) = V_0(0)$,

We set

- $\nu(t) = V_0(0)$,
- $\eta(t) = (\eta_\ell(t))_{\ell=1}^N$ with $\eta_\ell(t) = \sum_{k=1}^N Z_{\ell k} I_k(0)$.

We set

- $\nu(t) = V_0(0)$,
- $\eta(t) = (\eta_\ell(t))_{\ell=1}^N$ with $\eta_\ell(t) = \sum_{k=1}^N Z_{\ell k} I_k(0)$.

As an energy space, we define

$$\mathcal{H} = H \times H \times H \times \mathbb{C} \times \mathbb{C}^N,$$

where $H = \prod_{\ell=0}^N L^2(0, L)$.

We set

- $\nu(t) = V_0(0)$,
- $\eta(t) = (\eta_\ell(t))_{\ell=1}^N$ with $\eta_\ell(t) = \sum_{k=1}^N Z_{\ell k} I_k(0)$.

As an energy space, we define

$$\mathcal{H} = H \times H \times H \times \mathbb{C} \times \mathbb{C}^N,$$

where $H = \prod_{\ell=0}^N L^2(0, L)$.

The Hilbert space \mathcal{H} is equipped with the inner product:

We set

- $\nu(t) = V_0(0)$,
- $\eta(t) = (\eta_\ell(t))_{\ell=1}^N$ with $\eta_\ell(t) = \sum_{k=1}^N Z_{\ell k} I_k(0)$.

As an energy space, we define

$$\mathcal{H} = H \times H \times H \times \mathbb{C} \times \mathbb{C}^N,$$

where $H = \prod_{\ell=0}^N L^2(0, L)$.

The Hilbert space \mathcal{H} is equipped with the inner product:

$$\begin{aligned} \langle (V, I, W, \nu, \eta)^\top, (V^*, I^*, W^*, \nu^*, \eta^*)^\top \rangle_{\mathcal{H}} &= \sum_{\ell=0}^N \int_0^L (\theta_\ell V_\ell \overline{V_\ell^*} + \beta_\ell I_\ell \overline{I_\ell^*} + \gamma_\ell W_\ell \overline{W_\ell^*}) dx \\ &+ \delta Y \nu \overline{\nu^*} + \delta (\eta, (Z)^{-1} \eta^*)_{\mathbb{C}^N}, \quad \forall (V, I, W, \nu, \eta)^\top, (V^*, I^*, W^*, \nu^*, \eta^*)^\top \in \mathcal{H}. \end{aligned} \quad (6)$$

Where

$$\theta = (\theta_\ell)_{\ell=0}^N, \beta = (\beta_\ell)_{\ell=0}^N, \gamma = (\gamma_\ell)_{\ell=0}^N \in \prod_{\ell=0}^N L^\infty(0, L)$$

such that

We set

- $\nu(t) = V_0(0),$
- $\eta(t) = (\eta_\ell(t))_{\ell=1}^N$ with $\eta_\ell(t) = \sum_{k=1}^N Z_{\ell k} I_k(0).$

As an energy space, we define

$$\mathcal{H} = H \times H \times H \times \mathbb{C} \times \mathbb{C}^N,$$

where $H = \prod_{\ell=0}^N L^2(0, L).$

The Hilbert space \mathcal{H} is equipped with the inner product:

$$\begin{aligned} \langle (V, I, W, \nu, \eta)^\top, (V^*, I^*, W^*, \nu^*, \eta^*)^\top \rangle_{\mathcal{H}} &= \sum_{\ell=0}^N \int_0^L (\theta_\ell V_\ell \overline{V_\ell^*} + \beta_\ell I_\ell \overline{I_\ell^*} + \gamma_\ell W_\ell \overline{W_\ell^*}) dx \\ &+ \delta Y \nu \overline{\nu^*} + \delta (\eta, (Z)^{-1} \eta^*)_{\mathbb{C}^N}, \quad \forall (V, I, W, \nu, \eta)^\top, (V^*, I^*, W^*, \nu^*, \eta^*)^\top \in \mathcal{H}. \end{aligned} \quad (6)$$

Where

$$\theta = (\theta_\ell)_{\ell=0}^N, \beta = (\beta_\ell)_{\ell=0}^N, \gamma = (\gamma_\ell)_{\ell=0}^N \in \prod_{\ell=0}^N L^\infty(0, L)$$

such that

$$\theta_\ell \geq \mu_0, \beta_\ell \geq \mu_0, \gamma_\ell \geq \mu_0 \quad \text{a.e in } \Omega, \quad \ell = 0, \dots, N, \quad (7)$$

for some constant $\mu_0 > 0.$

The Operator \mathcal{A}

We define the operator $\mathcal{A} : D(\mathcal{A}) \rightarrow \mathcal{H}$ by :

The Operator \mathcal{A}

We define the operator $\mathcal{A} : D(\mathcal{A}) \rightarrow \mathcal{H}$ by :

$$D(\mathcal{A}) = \{(V, I, W, \nu, \eta) \in \mathcal{H}; V \in \prod_{\ell=0}^N H^1(0, L), I \in \prod_{\ell=0}^N H^1(0, L),$$
$$V_\ell(L) = \alpha_\ell I_\ell(L), \nu = V_0(0), \text{ and } \eta = (\eta_\ell)_{\ell=1}^N \text{ with } \eta_\ell = \sum_{k=1}^N Z_{\ell k} I_k(0)\}.$$
(8)

The Operator \mathcal{A}

We define the operator $\mathcal{A} : D(\mathcal{A}) \rightarrow \mathcal{H}$ by :

$$D(\mathcal{A}) = \left\{ (V, I, W, \nu, \eta) \in \mathcal{H}; V \in \prod_{\ell=0}^N H^1(0, L), I \in \prod_{\ell=0}^N H^1(0, L), \right. \\ \left. V_\ell(L) = \alpha_\ell I_\ell(L), \nu = V_0(0), \text{ and } \eta = (\eta_\ell)_{\ell=1}^N \text{ with } \eta_\ell = \sum_{k=1}^N Z_{\ell k} I_k(0) \right\}. \quad (8)$$

and

$$\mathcal{A}(V, I, W, \nu, \eta)^\perp = - \begin{pmatrix} (g_\ell V_\ell + a_\ell I_{\ell,x} + k_\ell W_\ell)_{\ell=0}^N \\ (r_\ell I_\ell + b_\ell V_{\ell,x})_{\ell=0}^N \\ (c_\ell W_\ell - V_\ell)_{\ell=0}^N \\ \frac{1}{\delta Y} \sum_{\ell=0}^N I_\ell(0) \\ \frac{1}{\delta} (V_\ell(0) - \nu)_{\ell=1}^N \end{pmatrix}.$$

Well posedness

If U is a sufficiently smooth solution of (2), then $U = (V, I, W, \nu, \eta) \in \mathcal{H}$ satisfies the first order evolution equation

$$\begin{cases} U_t = \mathcal{A}U, \\ U(0) = U_0. \end{cases} \quad (9)$$

Well posedness

If U is a sufficiently smooth solution of (2), then $U = (V, I, W, \nu, \eta) \in \mathcal{H}$ satisfies the first order evolution equation

$$\begin{cases} U_t = \mathcal{A}U, \\ U(0) = U_0. \end{cases} \quad (9)$$

THEOREM

There exists $\theta, \beta, \gamma \in \prod_{\ell=0}^N L^\infty(0, L)$ satisfying (7) such that the operator \mathcal{A} generates a C_0 semi-group of contractions on \mathcal{H} .

Using **Lumer-Phillip's Theorem** it is enough to prove that \mathcal{A} is m-dissipative.

Using **Lumer-Phillip's Theorem** it is enough to prove that \mathcal{A} is m-dissipative.

So we prove that

Using **Lumer-Phillip's Theorem** it is enough to prove that \mathcal{A} is m-dissipative.

So we prove that

$$\Re \langle \mathcal{A}U, U \rangle \leq -\frac{1}{2} \sum_{\ell=0}^N \int_{\Omega} (\theta_{\ell} g_{\ell} |V_{\ell}|^2 + 2\beta_{\ell} r_{\ell} |I_{\ell}|^2 + \gamma_{\ell} c_{\ell} |W_{\ell}|^2) dx - \sum_{\ell=0}^N \alpha_{\ell} |I_{\ell}(L)|^2 \leq 0, \quad (10)$$

Using **Lumer-Phillip's Theorem** it is enough to prove that \mathcal{A} is m-dissipative.

So we prove that

$$\Re \langle \mathcal{A}U, U \rangle \leq -\frac{1}{2} \sum_{\ell=0}^N \int_{\Omega} (\theta_{\ell} g_{\ell} |V_{\ell}|^2 + 2\beta_{\ell} r_{\ell} |I_{\ell}|^2 + \gamma_{\ell} c_{\ell} |W_{\ell}|^2) dx - \sum_{\ell=0}^N \alpha_{\ell} |I_{\ell}(L)|^2 \leq 0, \quad (10)$$

and

Using **Lumer-Phillip's Theorem** it is enough to prove that \mathcal{A} is m-dissipative.

So we prove that

$$\Re \langle \mathcal{A}U, U \rangle \leq -\frac{1}{2} \sum_{\ell=0}^N \int_{\Omega} (\theta_{\ell} g_{\ell} |V_{\ell}|^2 + 2\beta_{\ell} r_{\ell} |I_{\ell}|^2 + \gamma_{\ell} c_{\ell} |W_{\ell}|^2) dx - \sum_{\ell=0}^N \alpha_{\ell} |I_{\ell}(L)|^2 \leq 0, \quad (10)$$

and

$\lambda - \mathcal{A}$ is surjective for some $\lambda > 0$.

This implies that (2) is well-posed and its solution admits the following representation,

$$U(t) = e^{t\mathcal{A}} U_0, \quad t \geq 0.$$

Using **Lumer-Phillip's Theorem** it is enough to prove that \mathcal{A} is m-dissipative.

So we prove that

$$\Re \langle \mathcal{A}U, U \rangle \leq -\frac{1}{2} \sum_{\ell=0}^N \int_{\Omega} (\theta_{\ell} g_{\ell} |V_{\ell}|^2 + 2\beta_{\ell} r_{\ell} |I_{\ell}|^2 + \gamma_{\ell} c_{\ell} |W_{\ell}|^2) dx - \sum_{\ell=0}^N \alpha_{\ell} |I_{\ell}(L)|^2 \leq 0, \quad (10)$$

and

$\lambda - \mathcal{A}$ is surjective for some $\lambda > 0$.

This implies that (2) is well-posed and its solution admits the following representation,

$$U(t) = e^{t\mathcal{A}} U_0, \quad t \geq 0.$$

THEOREM

For any initial data $U_0 \in \mathcal{H}$, problem (9) admits a unique weak solution $U \in C^0([0, \infty); \mathcal{H})$.
Moreover if $U_0 \in D(\mathcal{A})$, problem (9) admits a unique strong solution $U \in C^1([0, \infty); \mathcal{H}) \cap C^0([0, \infty); D(\mathcal{A}))$.

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

The C_0 -semi group $(e^{tA})_{t \geq 0}$ is strongly stable on the energy space \mathcal{H} , i.e, for any $U_0 \in \mathcal{H}$, we have

$$\lim_{t \rightarrow \infty} \|e^{tA} U_0\|_{\mathcal{H}} = 0.$$

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

The C_0 -semi group $(e^{tA})_{t \geq 0}$ is strongly stable on the energy space \mathcal{H} , i.e, for any $U_0 \in \mathcal{H}$, we have

$$\lim_{t \rightarrow \infty} \|e^{tA} U_0\|_{\mathcal{H}} = 0.$$

The proof of the above theorem is based on the following lemmas :

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

The C_0 -semi group $(e^{tA})_{t \geq 0}$ is strongly stable on the energy space \mathcal{H} , i.e, for any $U_0 \in \mathcal{H}$, we have

$$\lim_{t \rightarrow \infty} \|e^{tA} U_0\|_{\mathcal{H}} = 0.$$

The proof of the above theorem is based on the following lemmas :

LEMMA (ARENDT-BATTY)

Let X be a reflexive Banach space and $(T(t))_{t \geq 0}$ be a bounded C_0 semi-group generated by A on X . If A has no eigenvalues on the imaginary axis and $\sigma(A) \cap i\mathbb{R}$ is countable, where $\sigma(A)$ denotes the spectrum of A . Then, $(T(t))_{t \geq 0}$ is stable.

To prove theorem (3), we need the following lemmas:

To prove theorem (3), we need the following lemmas:

LEMMA (1)

For all $\lambda \in \mathbb{R}$, we have

$$\ker(i\lambda - \mathcal{A}) = \{0\}.$$

To prove theorem (3), we need the following lemmas:

LEMMA (1)

For all $\lambda \in \mathbb{R}$, we have

$$\ker(i\lambda - \mathcal{A}) = \{0\}.$$

and

LEMMA (2)

For all $\lambda \in \mathbb{R}$, $i\lambda - \mathcal{A}$ is surjective.

To prove theorem (3), we need the following lemmas:

LEMMA (1)

For all $\lambda \in \mathbb{R}$, we have

$$\ker(i\lambda - \mathcal{A}) = \{0\}.$$

and

LEMMA (2)

For all $\lambda \in \mathbb{R}$, $i\lambda - \mathcal{A}$ is surjective.

Conclusion

- $(T(t))_{t \geq 1}$ is bounded,

To prove theorem (3), we need the following lemmas:

LEMMA (1)

For all $\lambda \in \mathbb{R}$, we have

$$\ker(i\lambda - \mathcal{A}) = \{0\}.$$

and

LEMMA (2)

For all $\lambda \in \mathbb{R}$, $i\lambda - \mathcal{A}$ is surjective.

Conclusion

- $(T(t))_{t \geq 1}$ is bounded,
- No eigenvalues on the imaginary axis, (Lemma (1))

To prove theorem (3), we need the following lemmas:

LEMMA (1)

For all $\lambda \in \mathbb{R}$, we have

$$\ker(i\lambda - \mathcal{A}) = \{0\}.$$

and

LEMMA (2)

For all $\lambda \in \mathbb{R}$, $i\lambda - \mathcal{A}$ is surjective.

Conclusion

- $(T(t))_{t \geq 1}$ is bounded,
- No eigenvalues on the imaginary axis, (Lemma (1))
- $\sigma(\mathcal{A}) \cap i\mathbb{R} = \emptyset$, (Lemma (1),(2))

To prove theorem (3), we need the following lemmas:

LEMMA (1)

For all $\lambda \in \mathbb{R}$, we have

$$\ker(i\lambda - \mathcal{A}) = \{0\}.$$

and

LEMMA (2)

For all $\lambda \in \mathbb{R}$, $i\lambda - \mathcal{A}$ is surjective.

Conclusion

- $(T(t))_{t \geq 1}$ is bounded,
- No eigenvalues on the imaginary axis, (Lemma (1))
- $\sigma(\mathcal{A}) \cap i\mathbb{R} = \emptyset$, (Lemma (1),(2))

Arendt -Batty theorem implies that $(e^{t\mathcal{A}})_{t \geq 0}$ is strongly stable.

Asymptotic Behavior of the eigenvalues of \mathcal{A} for $N = 1$ (2 cables)

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

In addition to (3), assume that the functions $\mathbf{a}, \mathbf{b}, \mathbf{c}, \mathbf{k}$ are constant vectors and that $\mathbf{r}_\ell = \mathbf{g}_\ell = \mathbf{0}, \forall \ell$. Then, there exists $n_0 \in \mathbb{N}^*$ such that the sequences of eigenvalues of $\mathcal{A}, (\lambda_n)_{n \geq n_0}$ and $(\lambda_m)_{m \geq n_0}$ satisfy the following asymptotic behavior:

Asymptotic Behavior of the eigenvalues of \mathcal{A} for $N = 1$ (2 cables)

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

In addition to (3), assume that the functions $\mathbf{a}, \mathbf{b}, \mathbf{c}, \mathbf{k}$ are constant vectors and that $\mathbf{r}_\ell = \mathbf{g}_\ell = \mathbf{0}, \forall \ell$. Then, there exists $n_0 \in \mathbb{N}^*$ such that the sequences of eigenvalues of \mathcal{A} , $(\lambda_n)_{n \geq n_0}$ and $(\lambda_m)_{m \geq n_0}$ satisfy the following asymptotic behavior:

$$\lambda_n = \begin{cases} \frac{\sqrt{a_0 b_0}}{2L} \log \left(\frac{1 - \frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}}}{1 + \frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}}} \right) + \frac{in\pi \sqrt{a_0 b_0}}{L} + o(1) & \text{if } \frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}} < 1, \\ or \\ \frac{\sqrt{a_0 b_0}}{2L} \log \left(\frac{\frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}} - 1}{\frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}} + 1} \right) + \frac{i(n + \frac{1}{2})\pi \sqrt{a_0 b_0}}{L} + o(1) & \text{if } \frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}} > 1, \end{cases} \quad (11)$$

and

Asymptotic Behavior of the eigenvalues of \mathcal{A} for $N = 1$ (2 cables)

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

In addition to (3), assume that the functions $\mathbf{a}, \mathbf{b}, \mathbf{c}, \mathbf{k}$ are constant vectors and that $\mathbf{r}_\ell = \mathbf{g}_\ell = \mathbf{0}, \forall \ell$. Then, there exists $n_0 \in \mathbb{N}^*$ such that the sequences of eigenvalues of \mathcal{A} , $(\lambda_n)_{n \geq n_0}$ and $(\lambda_m)_{m \geq n_0}$ satisfy the following asymptotic behavior:

$$\lambda_n = \begin{cases} \frac{\sqrt{a_0 b_0}}{2L} \log \left(\frac{1 - \frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}}}{1 + \frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}}} \right) + \frac{in\pi \sqrt{a_0 b_0}}{L} + o(1) & \text{if } \frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}} < 1, \\ \text{or} \\ \frac{\sqrt{a_0 b_0}}{2L} \log \left(\frac{\frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}} - 1}{\frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}} + 1} \right) + \frac{i(n + \frac{1}{2})\pi \sqrt{a_0 b_0}}{L} + o(1) & \text{if } \frac{\alpha_0 \sqrt{b_0}}{\sqrt{a_0}} > 1, \end{cases} \quad (11)$$

and

$$\lambda_m = \begin{cases} \frac{\sqrt{a_1 b_1}}{2L} \log \left(\frac{1 - \frac{\alpha_1 \sqrt{b_1}}{\sqrt{a_1}}}{1 + \frac{\alpha_1 \sqrt{b_1}}{\sqrt{a_1}}} \right) + \frac{i(m + \frac{1}{2})\pi \sqrt{a_1 b_1}}{L} + o(1) & \text{if } \frac{\alpha_1 \sqrt{b_1}}{\sqrt{a_1}} < 1, \\ \text{or} \\ \frac{\sqrt{a_1 b_1}}{2L} \log \left(\frac{\frac{\alpha_1 \sqrt{b_1}}{\sqrt{a_1}} - 1}{\frac{\alpha_1 \sqrt{b_1}}{\sqrt{a_1}} + 1} \right) + \frac{im\pi \sqrt{a_1 b_1}}{L} + o(1) & \text{if } \frac{\alpha_1 \sqrt{b_1}}{\sqrt{a_1}} > 1. \end{cases} \quad (12)$$

Exponential Energy decay for $N \geq 1$

We proved the exponential stability of our system for all $N \geq 1$, based on a frequency domain approach.

Exponential Energy decay for $N \geq 1$

We proved the exponential stability of our system for all $N \geq 1$, based on a frequency domain approach.

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

In addition to assumption (3), assume that $a_\ell, b_\ell \in W^{1,\infty}(0, L)$. Then, the semigroup $(e^{tA})_{t \geq 0}$ is exponentially stable, i.e, there exist constant M and $\epsilon > 0$ independent of U_0 such that

$$\|e^{tA}U_0\|_{\mathcal{H}} \leq Me^{-\epsilon t}\|U_0\|_{\mathcal{H}}, \quad t \geq 0. \quad (13)$$

Exponential Energy decay for $N \geq 1$

We proved the exponential stability of our system for all $N \geq 1$, based on a frequency domain approach.

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

In addition to assumption (3), assume that $a_\ell, b_\ell \in W^{1,\infty}(0, L)$. Then, the semigroup $(e^{tA})_{t \geq 0}$ is exponentially stable, i.e, there exist constant M and $\epsilon > 0$ independent of U_0 such that

$$\|e^{tA}U_0\|_{\mathcal{H}} \leq Me^{-\epsilon t}\|U_0\|_{\mathcal{H}}, \quad t \geq 0. \quad (13)$$

So, for this aim we use a theorem of Huang and Prüss, for which a C_0 -semi-group of contraction $(e^{tA})_{t \geq 0}$ is exponentially stable iff the following two conditions hold,

Exponential Energy decay for $N \geq 1$

We proved the exponential stability of our system for all $N \geq 1$, based on a frequency domain approach.

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

In addition to assumption (3), assume that $a_\ell, b_\ell \in W^{1,\infty}(0, L)$. Then, the semigroup $(e^{tA})_{t \geq 0}$ is exponentially stable, i.e, there exist constant M and $\epsilon > 0$ independent of U_0 such that

$$\|e^{tA}U_0\|_{\mathcal{H}} \leq Me^{-\epsilon t}\|U_0\|_{\mathcal{H}}, \quad t \geq 0. \quad (13)$$

So, for this aim we use a theorem of Huang and Prüss, for which a C_0 -semi-group of contraction $(e^{tA})_{t \geq 0}$ is exponentially stable iff the following two conditions hold,

$$(H1) \quad \rho(\mathcal{A}) \supset i\mathbb{R},$$

Exponential Energy decay for $N \geq 1$

We proved the exponential stability of our system for all $N \geq 1$, based on a frequency domain approach.

THEOREM (HAYEK-NICAISE-SALLOUM-WEHBE)

In addition to assumption (3), assume that $\mathbf{a}_\ell, \mathbf{b}_\ell \in W^{1,\infty}(0, L)$. Then, the semigroup $(e^{t\mathcal{A}})_{t \geq 0}$ is exponentially stable, i.e, there exist constant M and $\epsilon > 0$ independent of U_0 such that

$$\|e^{t\mathcal{A}}U_0\|_{\mathcal{H}} \leq Me^{-\epsilon t}\|U_0\|_{\mathcal{H}}, \quad t \geq 0. \quad (13)$$

So, for this aim we use a theorem of Huang and Prüss, for which a \mathcal{C}_0 -semi-group of contraction $(e^{t\mathcal{A}})_{t \geq 0}$ is exponentially stable iff the following two conditions hold,

$$(H1) \quad \rho(\mathcal{A}) \supset i\mathbb{R},$$

$$(H2) \quad \sup_{\lambda \in \mathbb{R}} \|(i\lambda - \mathcal{A})^{-1}\|_{L(\mathcal{H})} = O(1),$$

where, $\rho(\mathcal{A})$ denotes the resolvent of the operator \mathcal{A} .

1 INTRODUCTION

2 RESULTS ACHIEVED

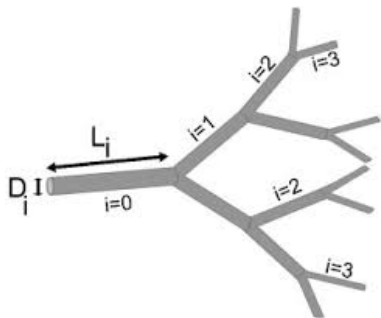
- Stabilization of a generalized Telegraph equation on a one dimensional star shaped network

3 CONTINUATION OF THE WORK

4 REFERENCES

- Generalization of the work to a tree-shaped network

- Generalization of the work to a tree-shaped network
- Later on, we can consider the generalized telegraph equation in an arbitrary graphs with general transmission conditions at the interior nodes.







1 INTRODUCTION

2 RESULTS ACHIEVED

- Stabilization of a generalized Telegraph equation on a one dimensional star shaped network

3 CONTINUATION OF THE WORK

4 REFERENCES

-  S. Nicaise. Stabilization and asymptotic behavior of a generalized telegraph equation. *Z. Angew. Math. Phys.*, 66(6):3221–3247, 2015.
-  S. Nicaise. Control and stabilization of 2×2 hyperbolic systems on graphs. *Math. Control Relat. Fields*, 7(1):53–72, 2017.
-  C. Zong and G. Q. Xu. Observability and controllability analysis of blood flow network. *Math. Control Relat. Fields*, 4(4):521–554, 2014.
-  F.L.Huang.Characteristics condition for exponential stability of linear dynamical systems in Hilbert spaces.*Ann.of Di.Eqs.*, 1:43 (56,1985).

THANK YOU FOR YOUR ATTENTION