# STABILIZATION OF COUPLED WAVE EQUATIONS WITH BOUNDARY DAMPING

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- Well-posedness and strong stability
- 4 Exponential stability (Case :  $\rho = 1$ )







3 Well-posedness and strong stability

4) Exponential stability (Case : ho = 1)

Polynomial stability (Case :  $\rho \neq 1$ )

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# **Motivation**

• In 1995, Cox-Zuazua studied the energy of the solution of the wave equation

$$\begin{cases} \rho^2(x)\partial_t^2 u(x,t) - \partial_x^2 u(x,t) = 0, & \text{in } ]0,1[\times]0, +\infty[\\ u(x,0) = u_0(x) , \partial_t u(x,0) = v_0(x), & \forall x \in [0,1] \end{cases}$$

under the damping condition

$$\partial_x u(1,t) + \partial_t u(1,t) = 0 \quad \forall t \in ]0, +\infty[.$$

• If the density  $\rho$  is suitably chosen, the dissipation of the energy through the right bound is sufficient to lead to exponential decay of solution.

 In the case of the density equal to one and no damping applied, i.e., when

$$u(1,t) = u(0,t) = 0 \ \forall t \in ]0, +\infty[,$$

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**Question** : How the stability properties are effected if we couple the exponentially stable wave equations to the conservative one ?

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3 Well-posedness and strong stability

### 4) Exponential stability (Case : $\rho = 1$ )

Polynomial stability (Case :  $\rho \neq 1$ )

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### Problem

• The functions  $u_1$  and  $u_2$  denote, respectively, the transverse displacement of a string of unit length in the case the density  $\rho^2 \neq 1$  and  $\rho = 1$ .

$$\begin{cases} \rho^2 \partial_t^2 u_1(x,t) - \partial_x^2 u_1(x,t) + \partial_t u_2(x,t) = 0, & \text{in } ]0,1[\times]0, +\infty[\\ \partial_t^2 u_2(x,t) - \partial_x^2 u_2(x,t) - \partial_t u_1(x,t) = 0, & \text{in } ]0,1[\times]0, +\infty[\\ u_1(0,t) = 0 & \forall t \in ]0, +\infty[\\ u_2(0,t) = u_2(1,t) = 0 & \forall t \in ]0, +\infty[ \end{cases}$$
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with the following initial conditions

 $u_1(x,0) = u_1^0(x), \ \partial_t u_1(x,0) = u_1^1(x), \ u_2(x,0) = u_2^0(x), \ \partial_t u_2(x,0) = u_2^1(x)$ 

and the boundary dissipation law

$$\partial_x u_1(1,t) + \partial_t u_1(1,t) = 0 \quad \forall t \in ]0, +\infty[.$$
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5/39 Problem

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5/39 Problem

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### Let us set

$$H_L^1(0,1) = \left\{ y \in H^1(0,1) : y(0) = 0 \right\}.$$

• Define the energy space H

 $H = H_L^1(0,1) \times L^2(0,1) \times H_0^1(0,1) \times L^2(0,1)$ 

equipped with the inner product defined by

$$(U, U_1)_H := \int_0^1 (u'\bar{u}'_1 + \rho^2 v \bar{v}_1 + y' \bar{y}'_1 + z\bar{z}_1) dx,$$
  
$$\forall U = (u, v, v, \bar{z}) \ U = (u, v, v, \bar{z}) \in H$$

6/39

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$$H^1_L(0,1) = \left\{ y \in H^1(0,1) : y(0) = 0 \right\}.$$

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• Let  $(u_1, u_2)$  be a regular solution of system (2.1)-(2.2), its associated total energy is defined by

$$E(t) = \frac{1}{2} \int_0^1 \left( |\partial_x u_1(x,t)|^2 + \rho^2 |\partial_t u_1(x,t)|^2 + |\partial_x u_2(x,t)|^2 + |\partial_t u_2(x,t)|^2 \right) dx$$
  
$$\forall t > 0.$$
  
(2.3)

We have the dissipation law

$$\frac{dE(t)}{dt} = -\left|\partial_x u_1(1,t)\right|^2.$$

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Hence, system (2.1)-(2.2) is dissipative in the sense that its associated energy is non increasing with respect to time.

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8/39 Problem

> **Question 1**: Does the energy decreases to zero? **Question 2**: When the energy approach to zero, how fast its decay is, and under what conditions?

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# Well-posedness

• Let  $U = (u_1, \partial_t u_1, u_2, \partial_t u_2)^t$ . Then the system (2.1)-(2.2) can be rewritten as an evolutionary equation in *H* 

$$\begin{cases} \partial_t U = AU(t), \quad t > 0 \\ U(0) = U_0 \end{cases}$$
(3.4)

where  $U_0 = (u_1^0, u_1^1, u_2^0, u_2^1)^t \in H$  is given.

•  $A: D(A) \longrightarrow H$  is defined by :

 $D(A) = \{(u, v, y, z) \in (H^2(0, 1) \cap H^1_L(0, 1)) \times H^1_L(0, 1) \times (H^2(0, 1) \cap H^1_0(0, 1)) \times H^1_0(0, 1)\}$ 

$$u'(1) + v(1) = 0\}$$

with

$$A(u,v,y,z)^{t} = \left(v,\frac{1}{\rho^{2}}\Delta u - \frac{1}{\rho^{2}}z,z,\Delta y + v\right)^{t}, \ \forall \ U = (u,v,y,z)^{t} \in D(A).$$

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The operator A is m-dissipative in the energy space H. In addition, the linear bounded operator  $A^{-1}$  is compact in H.

• Using Lumer-Phillips Theorem, the operator A generates a  $C_0$ -semigroup of contraction  $(e^{tA})_{t\geq 0}$  on H.

Theorem (Existence and uniqueness)

(1) If  $U_0 \in D(A)$ , then system (3.4) has a unique strong solution

 $U \in \mathcal{C}([0, +\infty[, D(A)) \cap \mathcal{C}^1([0, +\infty[, H)$ 

(2) If  $U_0 \in H$ , then system (3.4) has a unique weak solution

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$$U \in \mathcal{C}([0, +\infty[, D(A)) \cap \mathcal{C}^1([0, +\infty[, H)$$

(2) If  $U_0 \in H$ , then system (3.4) has a unique weak solution

 $U \in \mathcal{C}([0, +\infty[, H).$ 

# Strong stability

• *A* generates a contraction semigroup and its resolvent is compact in *H*, using Arendt-Batty theorem, system (2.1)-(2.2) is strongly stable, i.e,

$$\lim_{\to +\infty} ||e^{tA}U_0||_X = 0 \quad \forall U_0 \in H$$

if and only if A does not have pure imaginary eigenvalues.

### theorem (Strong Stability)

The semigroup of contraction  $(e^{tA})_{t\geq 0}$  is strongly stable on the energy space *H* if and only if the coefficient  $\rho$  satisfies the following condition

$$\rho^{4} + \left(2 - \frac{(k_{1}^{2} + k_{2}^{2})((k_{1}^{2} + k_{2}^{2})\pi^{2} - 1)}{k_{1}^{2}k_{2}^{2}\pi^{2}}\right)\rho^{2} + \frac{k_{1}^{2} + k_{2}^{2}}{k_{1}^{2}k_{2}^{2}\pi^{2}} + 1 \neq 0 \;\forall (k_{1}, k_{2}) \in K,$$
(3.5)

where  $K = \{(k_1, k_2) \in (\mathbb{N}^*)^2; k_1 \ge k_2 + 1 \text{ or } k_1 \le k_2 - 1\}.$ 

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#### Remark

### Let $\ensuremath{\mathcal{D}}$ be the set of solutions of

$$\rho^4 + \left(2 - \frac{(k_1^2 + k_2^2)((k_1^2 + k_2^2)\pi^2 - 1)}{k_1^2 k_2^2 \pi^2}\right)\rho^2 + \frac{k_1^2 + k_2^2}{k_1^2 k_2^2 \pi^2} + 1 = 0$$

then the condition (3.5) is equivalent to  $\rho^2 \notin \mathcal{D}$ .

Moreover,  $\mathcal{D}$  is a countable closed set which implies that the strong stability is true for almost any value of  $\rho^2$ . In addition,

$$\mathcal{D} \cap \mathbb{Q} = \emptyset. \tag{3.6}$$

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Well-posedness and strong stability

### 4 Exponential stability (Case : $\rho = 1$ )

Polynomial stability (Case :  $ho \neq 1$ )

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### Exponential stability (Case : $\rho = 1$ )

#### Theorem

Assume that  $\rho = 1$ , then the semigroup  $e^{tA}$  is exponentially stable, i.e., there exist a constant M > 0, and  $\omega > 0$  such that

$$|e^{tA}U_0|| \le Me^{-\omega t} ||U_0|| \quad \forall t > 0, U_0 \in D(A).$$
 (4.7)

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• Using Huang and Prüss, inequality (4.7) holds if and only if the following two conditions are satisfied :

(H1)  $i\mathbb{R} \subset \rho(A)$ ,

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$$\lim_{|\lambda|\to+\infty,\lambda\in\mathbb{R}} ||(i\lambda I - A)^{-1}|| < +\infty.$$

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# **Spectral Theory**

### Proposition

### In the case $\rho = 1$ , there exists $N \in \mathbb{N}$ such that

$$\{\mu_k\}_{k^* \in \mathbb{Z}, |k| \ge N} \subset \sigma(A) \tag{4.8}$$

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#### where

$$\mu_k = \frac{1}{2} \ln(\cos{(1)}) + ik\pi + o(1) \quad , k \in \mathbb{Z}^*, |k| \ge N$$

• Graphical interpretation of the exponential stability of the system.



**FIGURE:** Spectrum of *A*. Case :  $\rho = 1$ .

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# Let λ be an eigenvalue of A with associated eigenvector U = (u, v, y, z). AU = λU is equivalent to

$$\begin{cases} y'''' - 2\lambda^2 y'' + \lambda^2 (1 + \lambda^2) y = 0\\ -y''(1) - \frac{1}{\lambda} y'''(1) + \lambda y'(1) = 0\\ y''(0) = 0\\ y(1) = y(0) = 0 \end{cases}$$
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• The solution of (4.9) is given by :

$$y(x) = \sum_{i=1}^{4} C_i e^{t_i(\lambda)x}$$

where

$$\begin{split} t_1(\lambda) &= \sqrt{\lambda^2 + i\lambda}, \\ t_2(\lambda) &= -t_1(\lambda), \\ t_3(\lambda) &= \sqrt{\lambda^2 - i\lambda}, \\ t_4(\lambda) &= -t_3(\lambda). \end{split}$$

• The boundary conditions may be written as the following system :

$$M(\lambda)C(\lambda) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ e^{t_1} & e^{t_2} & e^{t_3} & e^{t_4} \\ t_1^2 & t_2^2 & t_3^2 & t_4^2 \\ h_\lambda(t_1)e^{t_1} & h_\lambda(t_2)e^{t_2} & h_\lambda(t_3)e^{t_3} & h_\lambda(t_4)e^{t_4} \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$
  
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• The asymptotic development of  $\frac{1}{\lambda^2}det(M(\lambda))$ , is given by

$$\tilde{f}(\lambda) = f_0(\lambda) + \frac{f_1(\lambda)}{\lambda} + O(\frac{1}{\lambda^2}),$$

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• Using the expansion of  $t_1$  and  $t_3$ :

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• The roots of  $f_0(\lambda)$ 

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 large enough.

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#### Remark

In order to obtain the best decay rate equal to the spectral abscissa

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we need to prove that the generalized eigenvectors associated to  $\mu_k$  form a Riesz basis of the energy space.

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## Condition (H2):

• Contradiction argument : Suppose that there exists a sequence  $(\lambda_n)_n \subset \mathbb{R}$  and a sequence  $(U^n)_n = (u^n, v^n, y^n, z^n)_n \subset D(A)$  verifying the following conditions

$$|\lambda_n| \longrightarrow +\infty,$$
 (4.11)

$$||U^n||_H = 1,$$
 (4.12)

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#### Step1 : (dissipation)

• Since  $(U^n)_n$  is uniformly bounded in H, and using the dissipation condition

$$|v^{n}(1)|^{2} = Re \langle (i\lambda_{n}I - A)U^{n}, U^{n} \rangle_{X} = o(1).$$
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Then

$$|u_x^n(1)| = o(1). (4.15)$$

and

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#### Step 2 :(Multiplier Method) we get $|y'(0)|^2 + |u'(0)|^2 - |y'(1)|^2 = o(1).$ (4.17)

• Under the condition y'(1) = o(1), one has

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**Step 3** :y'(1) = o(1). • Let  $Y = (u, u', y, y')^t$ , then (4.13) is equivalent to :

$$Y' = BY + G + \lambda F, \tag{4.18}$$

where

$$B = \begin{pmatrix} 0 & 1 & 0 & 0 \\ -\lambda^2 & 0 & i\lambda & 0 \\ 0 & 0 & 0 & 1 \\ -i\lambda & 0 & -\lambda^2 & 0 \end{pmatrix}, F = (F_j) = \begin{pmatrix} 0 \\ -if_1 \\ 0 \\ -if_3 \end{pmatrix},$$
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$$G = (G_j) = \begin{pmatrix} 0 \\ -f_2 - f_3 \\ 0 \\ -f_4 + f_1 \end{pmatrix}.$$

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$$Y(x) = e^{Bx}Y_0 + \int_0^x e^{B(x-z)}G(z)dz + \int_0^x \lambda e^{B(x-z)}F(z)dz$$
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where  $Y_0 = (u(0), u'(0), y(0), y'(0))^t = (0, u'(0), 0, y'(0))^t$ . • It easy to check that

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• When x = 1,

$$Y(1) = e^B Y_0 + o(\frac{1}{\lambda}).$$

• This is equivalent to

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• Using  $|y'(0)|^2 + |u'(0)|^2 - |y'(1)|^2 = o(1)$ , we get

$$(\cos(\frac{1}{2})\cos^2(\lambda) - \sin^2(\lambda)\sin^2(\frac{1}{2}) - 1)|y'(1)|^2 = o(1),$$

which is equivalent

$$\left(\cos^2(\lambda)(\cos^2(\frac{1}{2})-1)+\sin^2(\lambda)(-1-\sin^2(\frac{1}{2}))\right)|y'(1)|^2=o(1).$$

Since

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Well-posedness and strong stability

#### 4) Exponential stability (Case : ho = 1)

Polynomial stability (Case :  $\rho \neq 1$ )

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Polynomial stability (Case :  $\rho \neq 1$ )

## Polynomial stability (Case : $\rho \neq 1$ )

#### Theorem

Assume that  $\rho \neq 1$ . If  $\rho \in \mathbb{Q}$ , then there exists a constant C > 0 such that for every initial data  $U_0 \in D(A)$ , the energy of system (2.1)-(2.2) verifies the following estimate :

$$E(t) \le C \frac{1}{\sqrt{t}} ||U_0||^2_{D(A)}, \quad \forall t > 0.$$
 (5.24)

Using Borichev, Tomilov theorem a  $C_0$ -semigroup of contractions  $e^{tA}$  in a Hilbert space *H* satisfies (5.24) if and only if

(H1) 
$$i\mathbb{R} \subset \rho(A)$$
, and

(H2) 
$$\lim_{|\lambda| \to +\infty, \lambda \in \mathbb{R}} \sup_{|\lambda|^4} ||(i\lambda I - A)^{-1}|| < +\infty.$$

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 $\begin{array}{ll} \text{(H1)} & i\mathbb{R} \subset \rho(A), \\ & \text{and} \\ \text{(H2)} & \limsup_{|\lambda| \to +\infty, \lambda \in \mathbb{R}} \frac{1}{|\lambda|^4} ||(i\lambda I - A)^{-1}|| < +\infty. \end{array}$ 

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,  
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(H2)  $\lim_{x \to 0} \sup_{x \to 1} \frac{1}{2} ||(i)|_{x \to 1}$ 

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$$\limsup_{|\lambda| \to +\infty, \lambda \in \mathbb{R}} \frac{|\lambda|^4}{|\lambda|^4} ||(i\lambda I - A)^{-1}|| < +\infty.$$

### **Spectral Theory**

#### Proposition

Assume that  $\rho > 1$  . There exists  $N \in \mathbb{N}^*$  such that

$$\{\lambda_k\}_{k\in\mathbb{Z}^*,|k|\ge N}\cup\{\mu_k\}_{k\in\mathbb{Z},|k|\ge N}\subset\sigma(A)$$
(5.25)

#### where

$$\lambda_k = ik\pi + o(1), k \in \mathbb{Z}^*, |k| \ge N,$$
  
$$\mu_k = \frac{-1}{2\rho} \ln\left(\frac{\rho+1}{\rho-1}\right) + i\frac{\pi k}{\rho} + o(1), k \in \mathbb{Z}, |k| \ge N$$

Moreover for all  $|k| \ge N$ , the eigenvalues  $\lambda_k$  and  $\mu_k$  are simple.

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#### • Graphical interpretation of the polynomial stability of the system



## FIGURE: Spectrum of A. Case : $\rho = 2$ .

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# Let λ be an eigenvalue of A with associated eigenvector U = (u, v, y, z). AU = λU is equivalent to

$$\begin{cases} y''' - \lambda^2 (1 + \rho^2) y'' + \lambda^2 (1 + \lambda^2 \rho^2) y = 0 \\ -y''(1) - \frac{1}{\lambda} y'''(1) + \lambda y'(1) = 0 \\ y''(0) = 0 \\ y(1) = y(0) = 0. \end{cases}$$
(5.26)

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AU = λU is equivalent to

$$\begin{cases} y'''' - \lambda^2 (1 + \rho^2) y'' + \lambda^2 (1 + \lambda^2 \rho^2) y = 0\\ -y''(1) - \frac{1}{\lambda} y'''(1) + \lambda y'(1) = 0\\ y''(0) = 0\\ y(1) = y(0) = 0. \end{cases}$$
(5.26)

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• $\lambda$  is an eigenvalue of *A* if and only if there is a non trivial solution of (5.26) which satisfies the boundary conditions.

• The general solution of (5.26) is given by

$$y(x) = \sum_{i=1}^{4} C_i e^{t_i(\lambda)x}$$

#### where

$$t_1(\lambda) = \sqrt{\frac{\lambda^2(1+\rho^2) + \lambda\sqrt{\lambda^2(1-\rho^2)^2 - 4}}{2}}, \quad t_2(\lambda) = -t_1(\lambda)$$
  
$$t_3(\lambda) = \sqrt{\frac{\lambda^2(1+\rho^2) - \lambda\sqrt{\lambda^2(1-\rho^2)^2 - 4}}{2}}, \quad t_4(\lambda) = -t_3(\lambda).$$

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• The boundary conditions may be written as the following system :

$$M(\lambda)C(\lambda) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ e^{t_1} & e^{t_2} & e^{t_3} & e^{t_4} \\ t_1^2 & t_2^2 & t_3^2 & t_4^2 \\ h_\lambda(t_1)e^{t_1} & h_\lambda(t_2)e^{t_2} & h_\lambda(t_3)e^{t_3} & h_\lambda(t_4)e^{t_4} \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

where  $h_{\lambda}(t) = t^2 + \frac{1}{\lambda}t^3 - \lambda t$ .

• A non trivial solution y exists if and only if the determinant of  $M(\lambda)$  vanishes.

• Let  $f(\lambda) = det M(\lambda)$ , thus the characteristic equation

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• As we did before, using the asymptotic expansions of  $t_1$  and  $t_3$  and multipling the third and the fourth line of  $M(\lambda)$  by  $\frac{1}{\lambda^2}$ , we obtain the following asymptotic development

$$f(\lambda) = f_0(\lambda) + \frac{f_1(\lambda)}{\lambda^2} + \frac{f_2(\lambda)}{\lambda^4} + O(\frac{1}{\lambda^4}).$$

• Using the Rouché's theorem, and for  $\lambda$  large enough, the roots of f are close to those of  $f_0$  where

$$f_0(\lambda) = (\rho - 1)^2 (\rho + 1)^2 e^{t_1 + t_3} (1 - e^{-2t_3}) (\rho + 1 + (\rho - 1)e^{-2t_1}).$$

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• $f_0$  has two families of roots that we denote  $\lambda_k^0$  and  $\mu_k^0$ . First case :

 $e^{-2t_3}=1 \iff 2t_3=2ik\pi, \ k\in\mathbb{Z}.$ 

Since  $t_3 \sim \lambda$ , we get the first family of roots of  $f_0$  is

 $\lambda_k^0 \sim ik\pi + o(1), \;\; k \in \mathbb{Z} \;\;$ large enough.

Second case :  $e^{-2t_1} = \frac{\rho+1}{\rho-1}$ , then

$$-2t_1 = \ln\left(\frac{\rho+1}{\rho-1}\right) + ik\pi, \ k \in \mathbb{Z}.$$

The second family of roots of  $f_0$  is

$$\mu_k^0 = -\frac{1}{2\rho} \ln\left(\frac{\rho+1}{\rho-1}\right) + i\frac{k\pi}{\rho} + o(1), \quad k \in \mathbb{Z}.$$

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#### Conclusion :

• There exists a sequence  $(\lambda_n)_n$  of roots of f such that

$$\lambda_k = ik\pi + o(1)$$
 as  $k \longrightarrow +\infty$ .

#### Remark

In order to obtain the optimal energy decay rate we should investigate the asymptotic behavior for the first family of eigenvalues near the imaginary axis.

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## Condition (H2)

• Using again an argument of contradiction : there exists a sequence  $(\lambda_n)_n \subset \mathbb{R}$  and a sequence  $(U^n)_n = (u^n, v^n, y^n, z^n)_n \subset D(A)$  verifying the following conditions

$$|\lambda_n| \longrightarrow +\infty,$$
 (5.27)

$$||U^n||_H = 1, (5.28)$$

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$$\lambda_n^4(i\lambda_n I - A)(u^n, v^n, y^n, z^n) = (f_1^n, f_2^n, f_3^n, f_4^n) \longrightarrow 0 \text{ in } H.$$
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• Our aim is to prove that  $||U^n||_H = o(1)$ . This contradicts equation (5.28).

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$$y_x(1) = o(1).$$
 (5.30)

 To simplify the computation we translate the problem and we prove that

$$y_x(0)=o(1).$$

• If this is not true, then there exists a constant c such that

$$y_x^n(0) \ge c, \quad \forall n \in \mathbb{N}.$$

• (5.29) can be written as

$$U' = BU + F$$

where

$$U = \begin{pmatrix} u \\ u_x \\ y \\ y_x \end{pmatrix}, B = \begin{pmatrix} 0 & 1 & 0 & 0 \\ -\rho^2 \lambda^2 & 0 & i\lambda & 0 \\ 0 & 0 & 0 & 1 \\ -i\lambda & 0 & -\lambda^2 & 0 \end{pmatrix} \text{ and } F = \begin{pmatrix} 0 \\ f \\ 0 \\ g \end{pmatrix}.$$

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• Applying the boundary conditions u(1) = y(1) = 0, and using the fact that  $y'(0) \ge c$ , it follows that there exist  $m, k \in \mathbb{Z}$  with the same parity such that

$$\begin{cases} \lambda - \frac{1}{2\lambda(\rho^2 - 1)} = m\pi + \frac{O(1)}{\lambda^2}, \\ \lambda \rho - \frac{1}{2\lambda\rho(\rho^2 - 1)} = k\pi + \frac{O(1)}{\lambda}. \end{cases}$$
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$$\begin{cases} u = \sin(\lambda - \frac{1}{2\lambda(\rho^2 - 1)})y'(0) + \frac{o(1)}{\lambda^3}, \\ y = \sin(\lambda\rho - \frac{1}{2\lambda\rho(\rho^2 - 1)})y'(0) + \frac{o(1)}{\lambda^4}. \end{cases}$$
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• Applying the boundary conditions u(1) = y(1) = 0, and using the fact that  $y'(0) \ge c$ , it follows that there exist  $m, k \in \mathbb{Z}$  with the same parity such that

$$\begin{cases} \lambda - \frac{1}{2\lambda(\rho^2 - 1)} = m\pi + \frac{O(1)}{\lambda^2}, \\ \lambda \rho - \frac{1}{2\lambda\rho(\rho^2 - 1)} = k\pi + \frac{O(1)}{\lambda}. \end{cases}$$
(5.32)

$$\begin{cases} \lambda^{2} = m^{2}\pi^{2} + \frac{\pi}{\rho^{2} - 1} + \frac{O(1)}{\lambda}, \\ \lambda^{2}\rho^{2} = k^{2}\pi^{2} - \frac{\pi}{\rho(\rho^{2} - 1)} + o(1). \end{cases}$$
(5.33)

$$\rho^2 m^2 - k^2 = -\frac{\rho^3 + 1}{\pi \rho (\rho^2 - 1)} + o(1).$$

• Let  $c = -\frac{\rho^3 + 1}{\pi \rho(\rho^2 - 1)}$ . • Since we have assumed that  $\rho = \frac{p}{q}$  for some  $(p, q) \in \mathbb{N}^*$ , we deduce

$$\frac{pm-kq}{q} = \frac{c}{pm+kq} + \frac{o(1)}{pm+kq}.$$

*i*) If pm - qk = 0 for an infinity number of pairs (m, k), then c = o(1) and this a contradiction.

*ii*) Else  $pm - kn \neq 0$  for  $\lambda$  large enough and then

$$\frac{1}{q} \le |\frac{c}{pm + kq}| + |\frac{o(1)}{pm + kq}| \le o(1),$$

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